

# New opportunities offered by heavy ion transfer reactions

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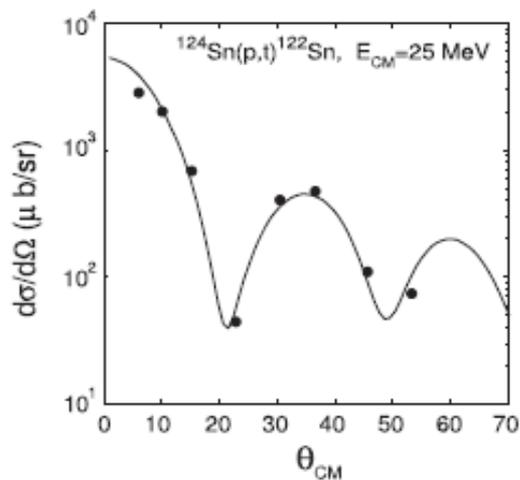
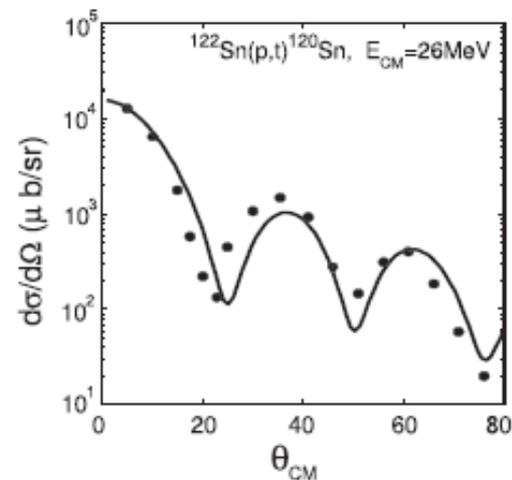
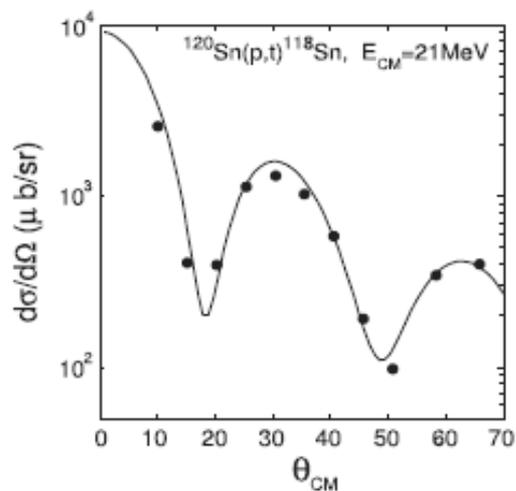
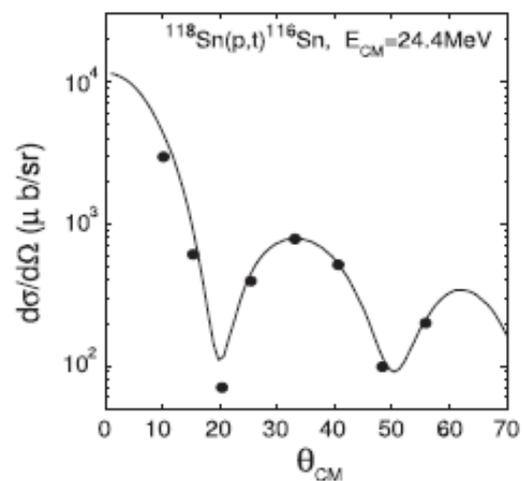


NUSPRASEN Workshop on Nuclear Reactions,  
Warsaw (Poland) - January 22-24, 2018

**Part I**  
**nucleon nucleon correlations**

**Part II**  
**production of neutron-rich nuclei**

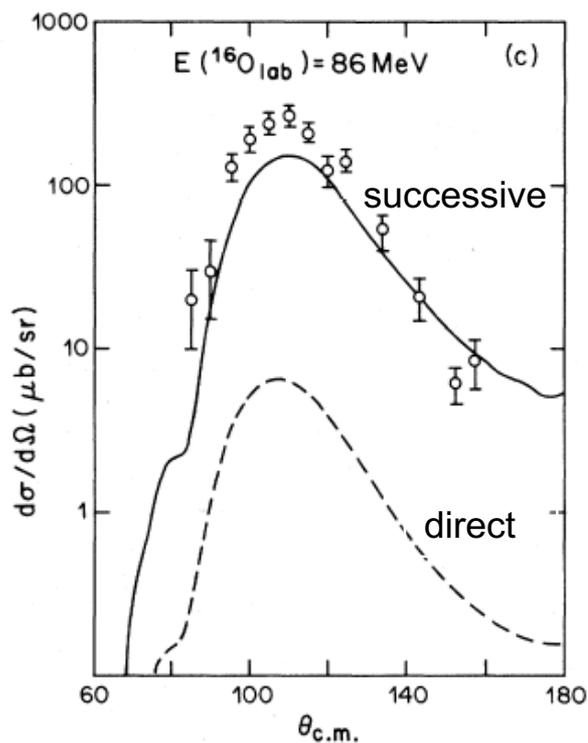
## (p,t) reactions : absolute cross sections



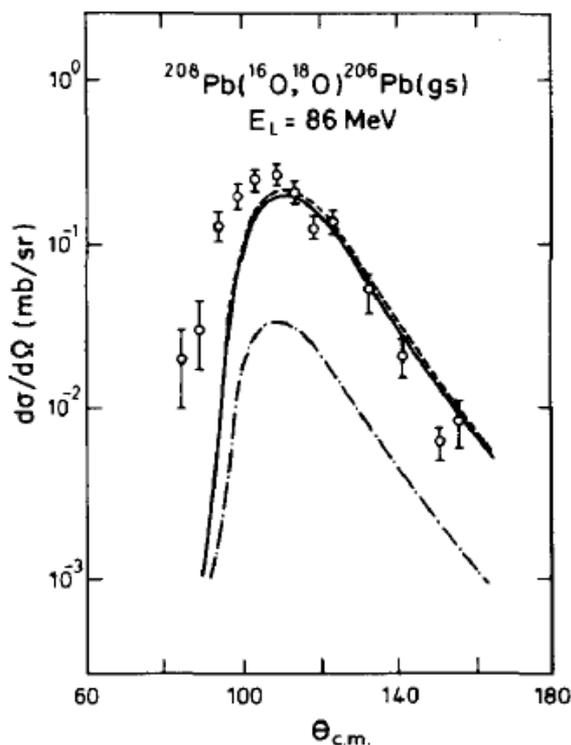
recently performed  
calculations for two  
neutron transfer  
reactions match the  
experimental data  
with high accuracy

# Absolute cross sections for one and two-nucleon transfer reactions

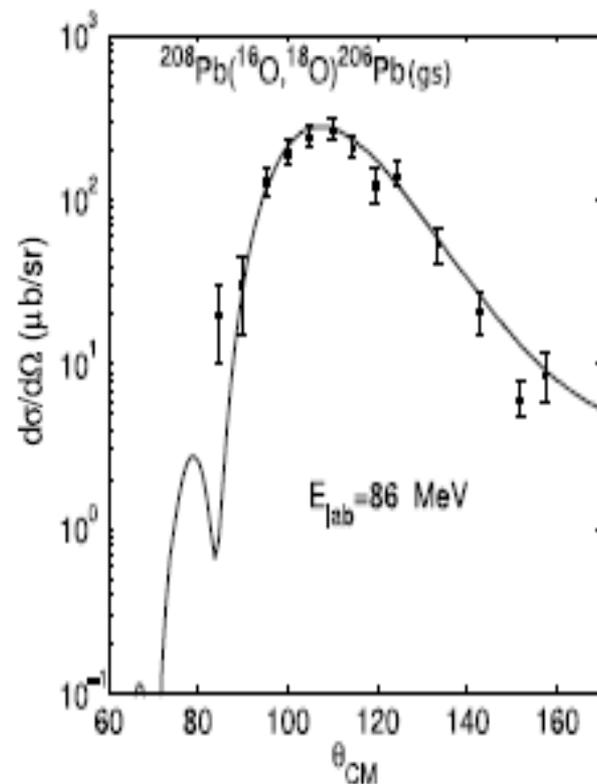
## $^{208}\text{Pb}(^{16}\text{O}, ^{18}\text{O}_{\text{g.s.}})^{206}\text{Pb}$



B.F.Bayman and J.Chen,  
PRC26(1982)1509



E.Maglione, G.Pollarolo, A.Vitturi,  
R.A.Brogia and A.Winther  
PLB162(1985)59



G.Potel, A.Idini, F.Barranco,  
E.Vigezzi and R.A.Brogia  
Rep.Prog.Phys.76(2013)106301

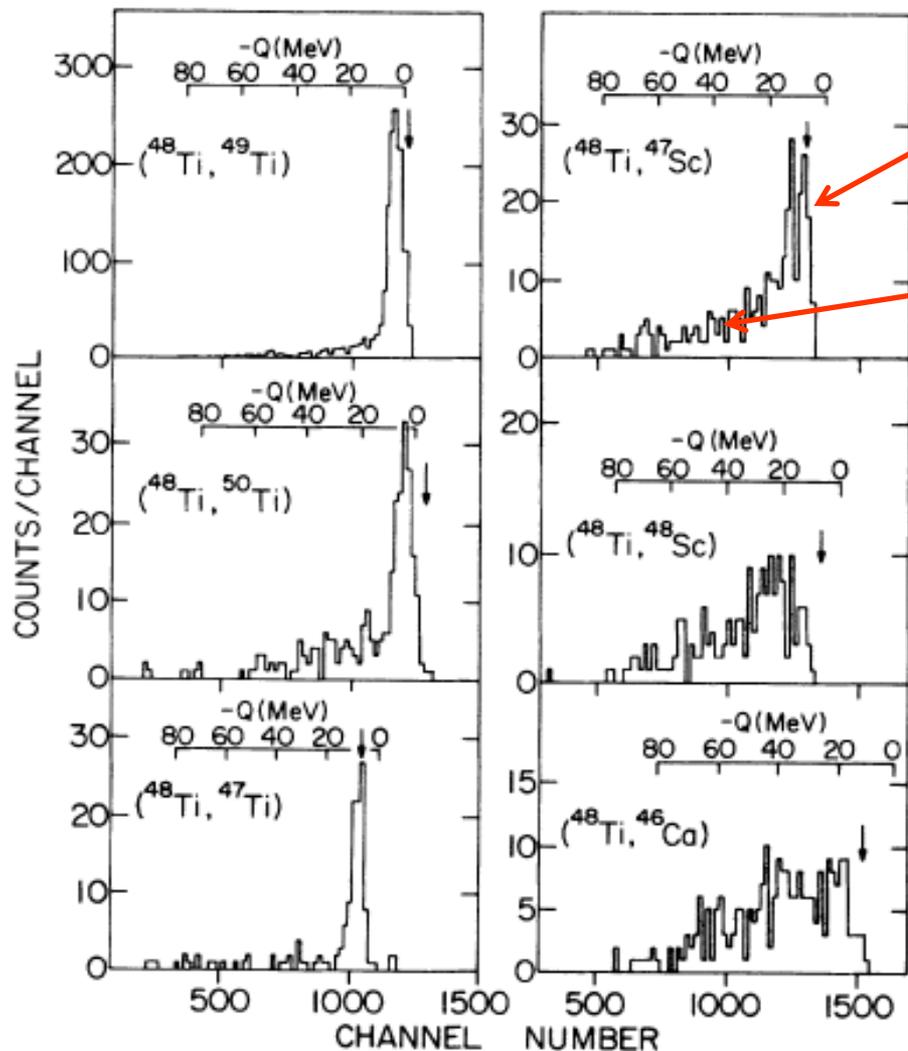
full quantum mechanical

semiclassical

full quantum mechanical

# Quasi-elastic (QE) and Deep Inelastic (DIC) regime in multinucleon transfer reactions : Q-values

$^{48}\text{Ti} + ^{208}\text{Pb}$ ,  $E_{\text{lab}} = 300 \text{ MeV}$ ,  $\theta_{\text{lab}} = 55^\circ$



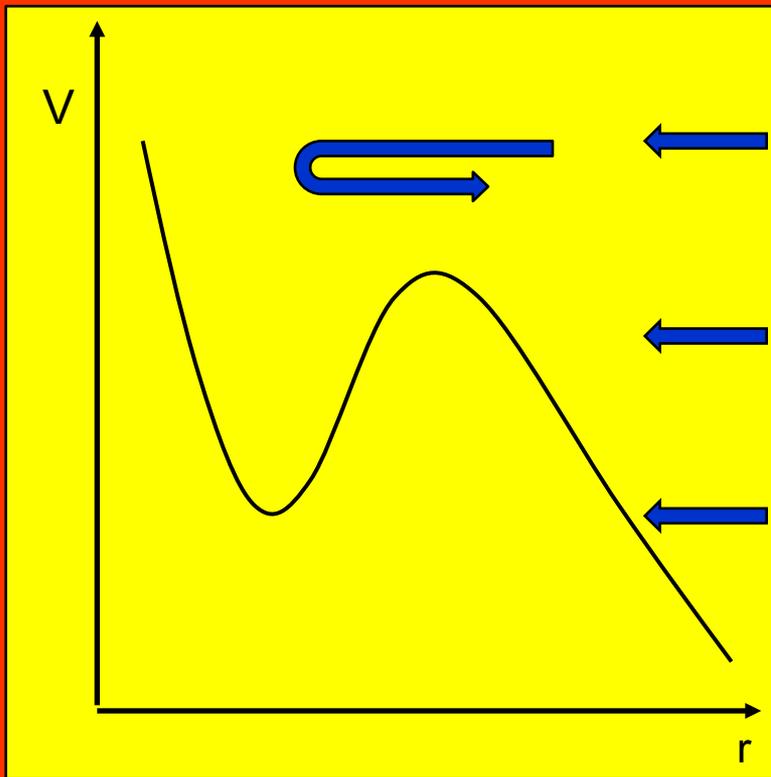
QE

DIC

with heavy ions the general difficulty met in both experiments and theory is to deal with limited energy resolutions and with the presence of both QE and DIC components

**Transfer studies at energies below the Coulomb barrier: advantages**

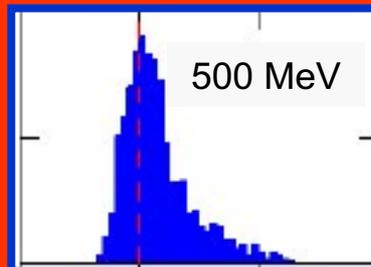
$^{116}\text{Sn}(^{60}\text{Ni}, ^{62}\text{Ni})$   
 $Q_{\text{gs}} = +1.3 \text{ MeV}$



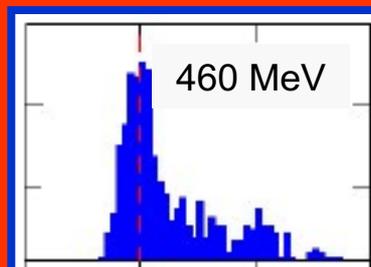
$E > E_b$

$E \sim E_b$

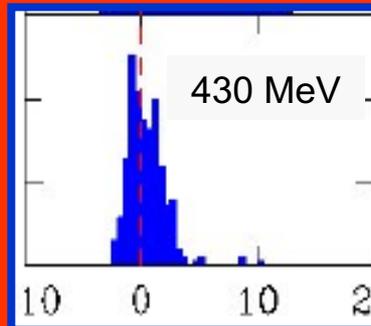
$E < E_b$



large number of open channels, DIC components



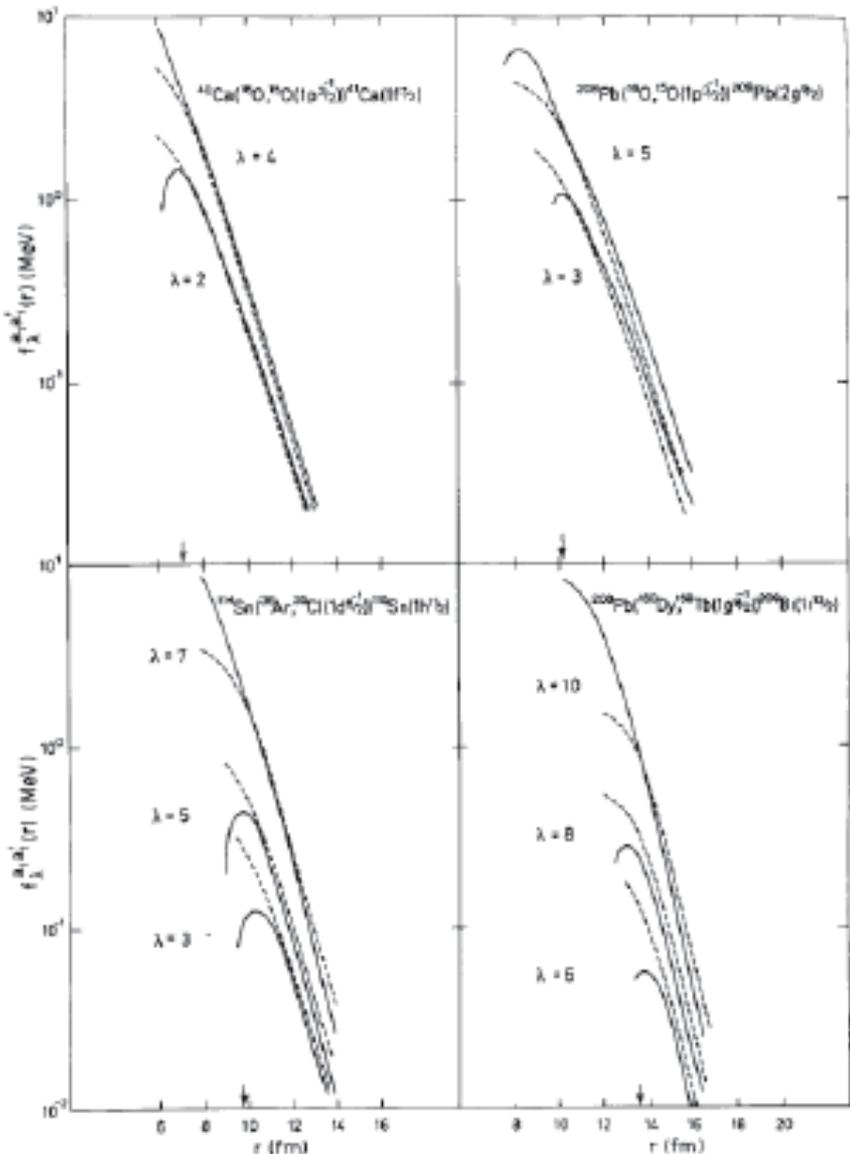
the number of open channels and uncertainties with the nuclear potential reduce



one gets very narrow Q-value distributions: no evaporation effects

**below the barrier one can probe nucleon-nucleon correlations as close as possible to the ground to ground states**

# Quasi elastic processes : form factors



$$\langle \omega_{\beta} | (V_{\gamma} - U_{\gamma}) | \psi_{\gamma} \rangle = f_{\beta\gamma}(\vec{k}, \vec{r})$$

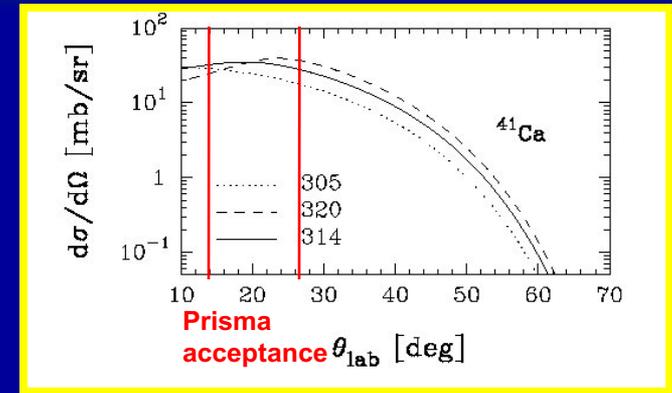
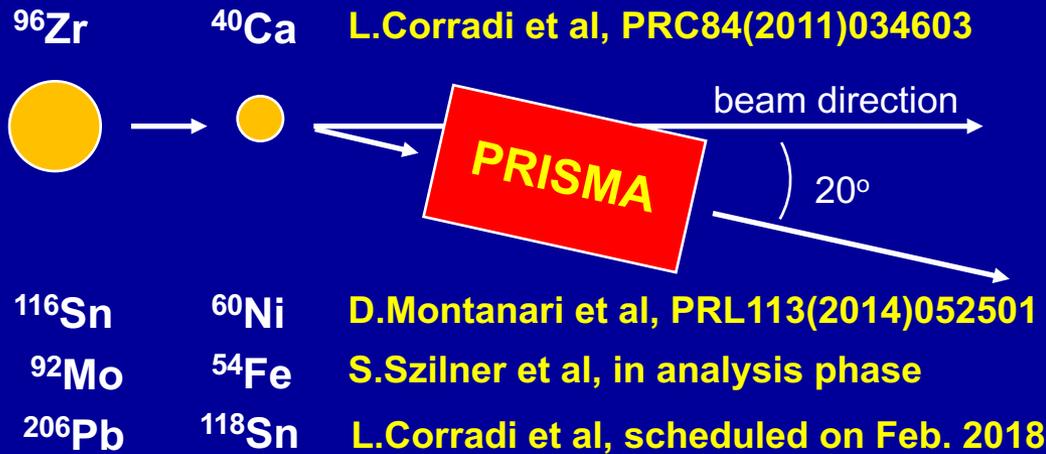
$$f_{\beta\gamma}(\vec{k}, \vec{r}) \sim e^{i\sigma_{\beta\gamma}t} f_{\beta\gamma}(0, \vec{r})$$

the form factor is a matrix element between initial and final states in the transfer process and reflects nuclear structure properties of the donor and acceptor binary partners

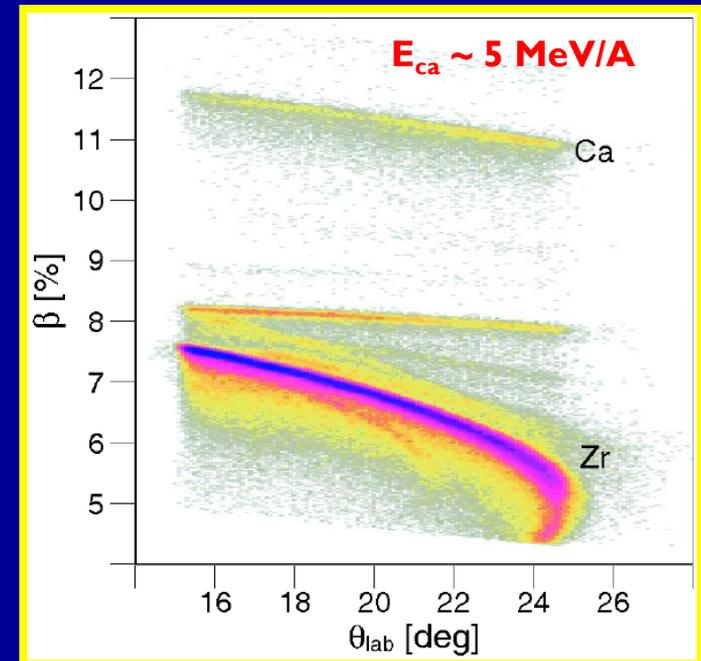
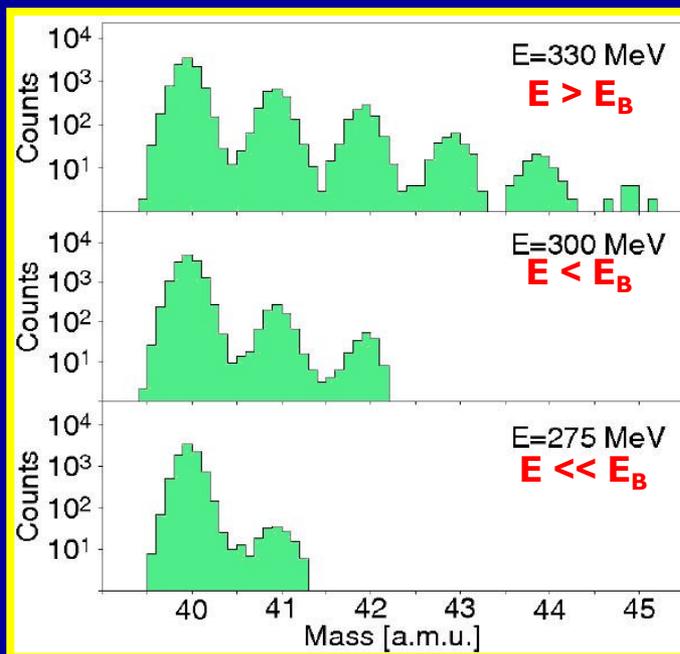
$$f_{\beta\gamma}(0, r) \propto \frac{1}{K_{a_1} r} e^{-K_{a_1} r}$$

the form factor has an exponential shape in its tail region

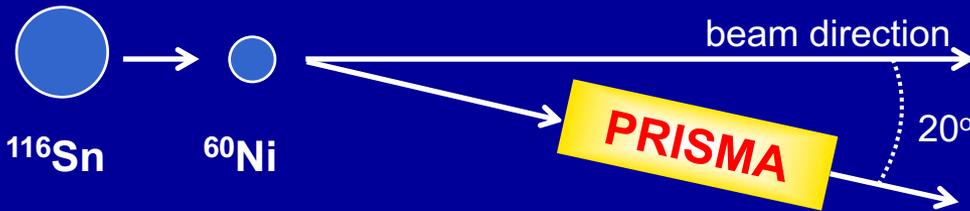
# Detection of (light) target like ions in inverse kinematics with PRISMA



MNT channels have been measured down to 25 % below the Coulomb barrier



# Detection of (light) target like ions in inverse kinematics with PRISMA



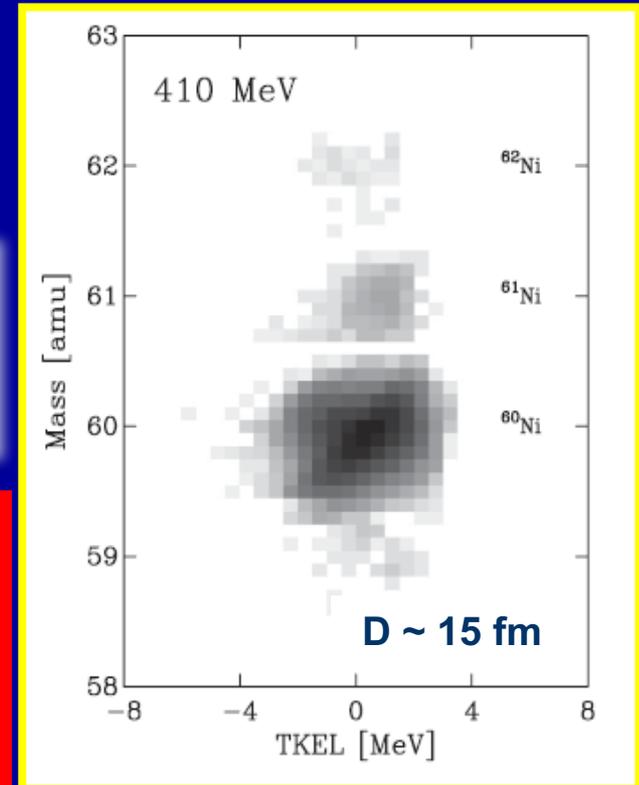
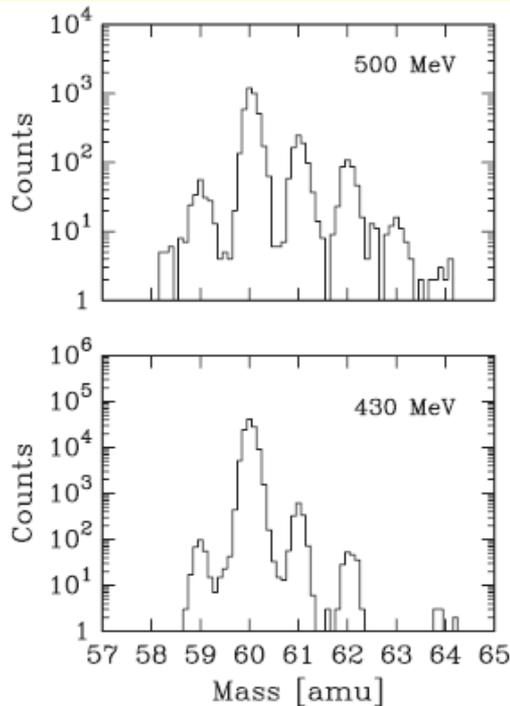
**excellent mass resolution  
at energies below the  
Coulomb barrier**

**Excitation function**

$$E_{\text{beam}} = 410 - 500 \text{ MeV}$$

$$(D \sim 12.3 \text{ to } 15.0 \text{ fm})$$

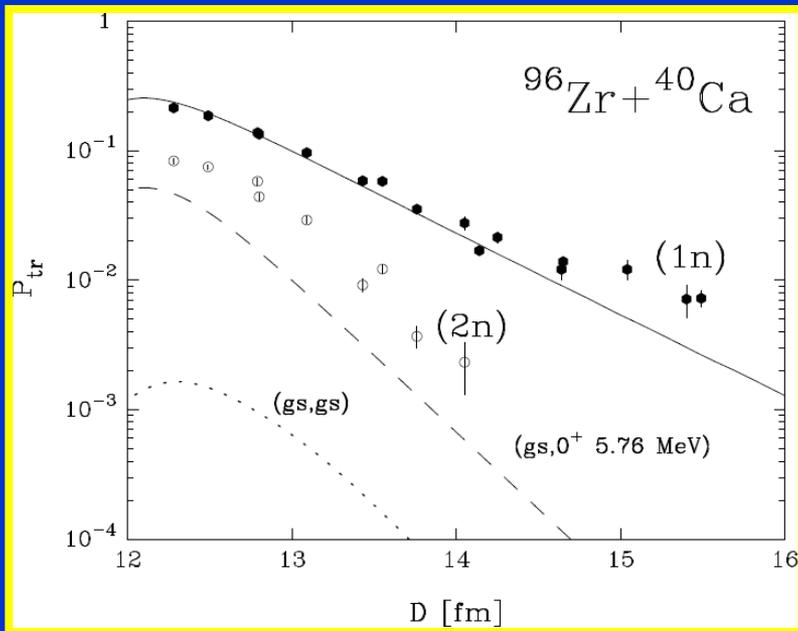
**ground to ground  
state Q-values  
very close to  
optimum Q-values  
(~ 0 MeV)**



g.s. Q-values	+1n	+2n	+3n	+4n
$^{96}\text{Zr} + ^{40}\text{Ca}$	+ 0.51	+ 5.53	+ 5.24	+ 9.64
$^{116}\text{Sn} + ^{60}\text{Ni}$	- 1.74	+ 1.31	- 2.15	- 0.24

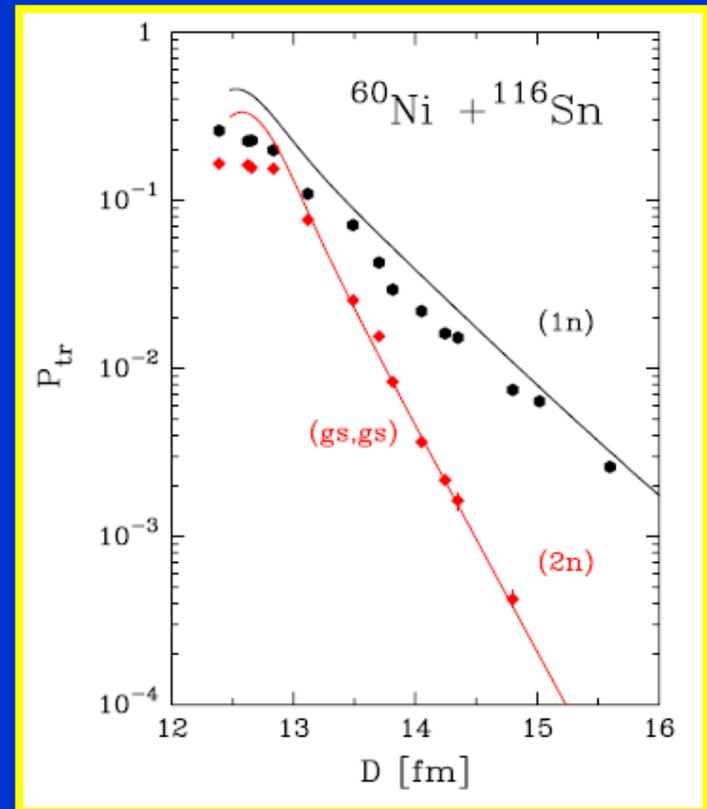
# Transfer probabilities : comparison between exp and microscopic theory

$Q_{g.s.}$  for +2n + 5.5 MeV , far from  $Q_{opt}$  ( $\sim 0$  MeV)



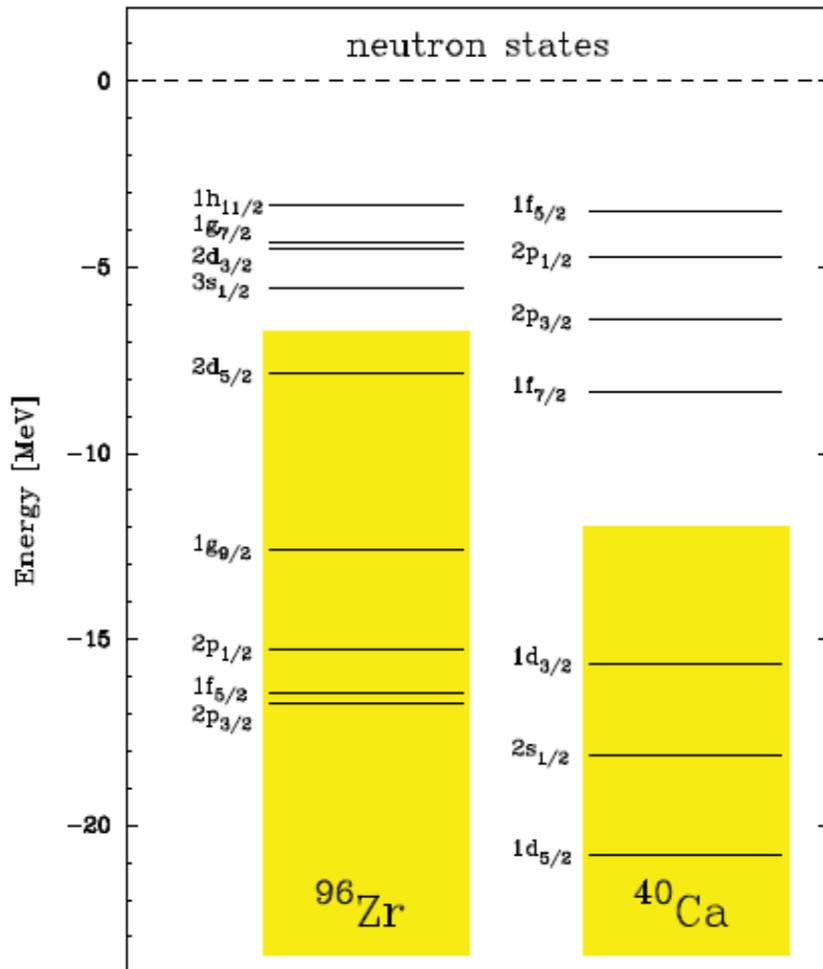
L.Corradi, S.Szilner, G.Pollarolo et al,  
PRC84(2011)034603

$Q_{g.s.}$  for +2n very close to  $Q_{opt}$  ( $\sim 0$  MeV)



D.Montanari, L.Corradi, S.Szilner,  
G.Pollarolo et al, PRL113(2014)052501

# One particle transfer (semiclassical theory)



to obtain the total transfer probability we summed over all possible transitions that can be constructed from the single particle states in projectile and target

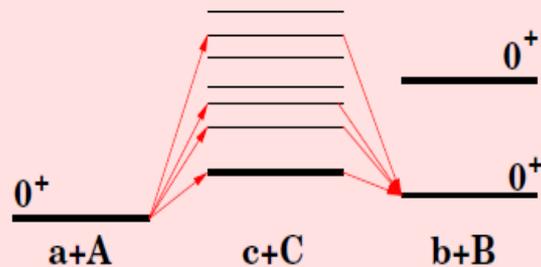
the set of single particle states covers a full shell below the Fermi level for  $^{96}\text{Zr}$  and a full shell above for  $^{40}\text{Ca}$

$$c_{\beta}(\ell) = \frac{1}{i\hbar} \int_{-\infty}^{+\infty} \langle \psi_{\beta} | (V_{\alpha} - U_{\alpha}) | \psi_{\alpha} \rangle_{\mathcal{R}} e^{i(E_{\beta} - E_{\alpha})t/\hbar} dt$$

$$P_{\beta}(\ell) = P_{(a_1, a'_1)}(\ell) = \sum_{m'_1, m_1} |c_{\beta}(\ell)|^2$$

## Two particle transfer (semiclassical theory, microscopic calculations)

$$c_{\beta}(\ell) = c_{\beta}^{(1)} + c_{\beta}^{ort} + c_{\beta}^{succ}$$



**3 terms : simultaneous, orthogonal and successive**

$$c_{\beta}(\ell) = (c_{\beta})_{(1)} + (c_{\beta})_{ort} + (c_{\beta})_{succ}$$

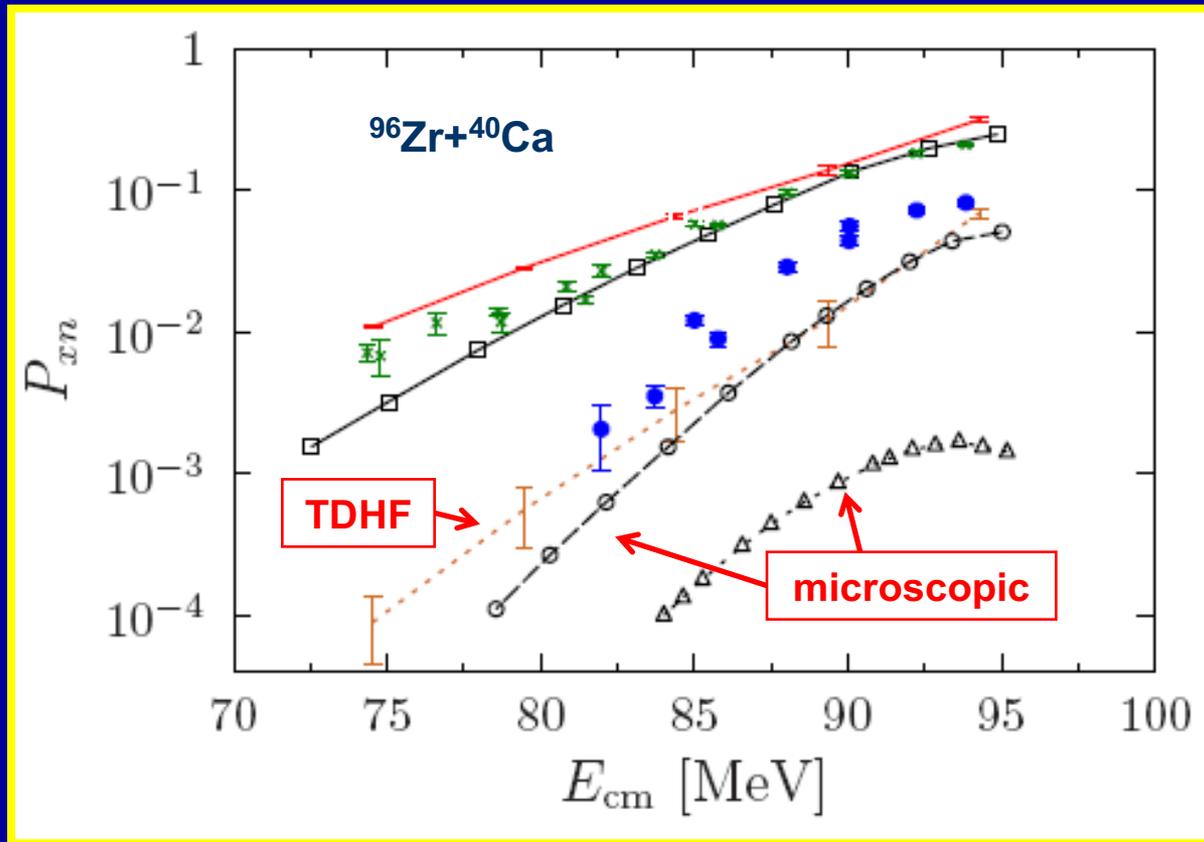
**only the successive term contributes to the transfer amplitude**

**only the  $0^+$  to  $0^+$  transitions can be reduced to a simple expression**

$$\begin{aligned} (c_{\beta})_{succ} &= \frac{1}{\hbar^2} \sum_{a_1, a'_1} B^{(A)}(a_1 a_1; 0) B^{(a)}(a'_1 a'_1; 0) 2 \frac{(-1)^{j_1 + j'_1}}{\sqrt{(2j_1 + 1)} \sqrt{(2j'_1 + 1)}} \sum_{m_1 m'_1} (-1)^{m_1 + m'_1} \\ &\times \int_{-\infty}^{+\infty} dt f_{m_1 m'_1}(\mathcal{R}) e^{i[(E_{\beta} - E_{\gamma})t + \delta_{\beta\gamma}(t) + \hbar(m'_1 - m_1)\Phi(t)]/\hbar} \\ &\times \int_{-\infty}^t dt f_{-m_1 - m'_1}(\mathcal{R}) e^{i[(E_{\gamma} - E_{\alpha})t + \delta_{\gamma\alpha}(t) - \hbar(m'_1 - m_1)\Phi(t)]/\hbar} . \end{aligned}$$

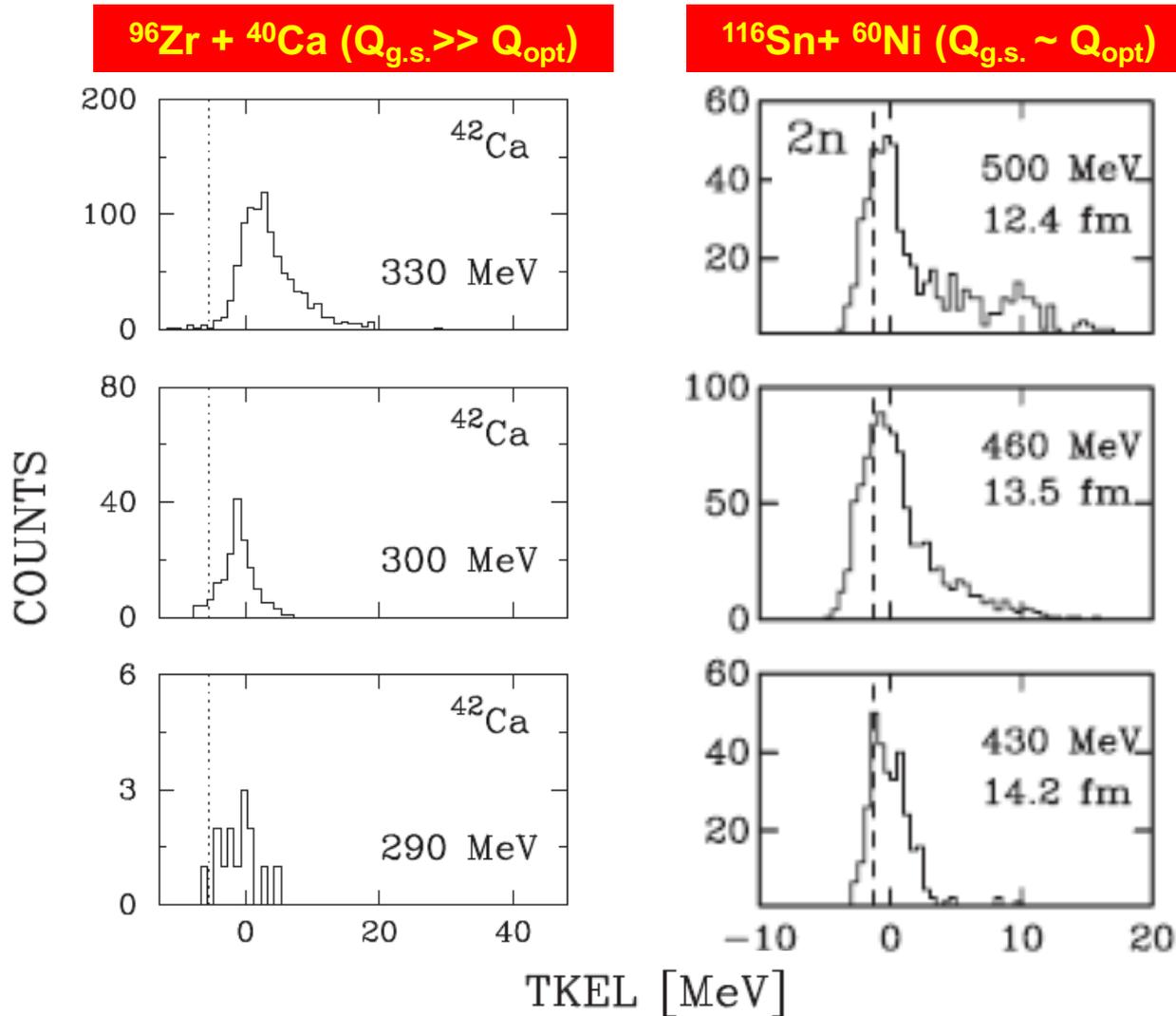
# Sub-barrier transfer : comparison between data and TDHF calculations

Data : Legnaro Nat. Lab.  
Theory: simplified TDHF+BCS limit



G.Scamps and D.Lacroix,  
EPJ Web Conf. 86, (2015) 00042

# Total kinetic energy loss distributions for the +2n channels



The key difference between the two cases lies in the relative position of the g.s. Q-values with respect to the optimum Q-values for neutrons

# Probing directly the population to the ground to ground states



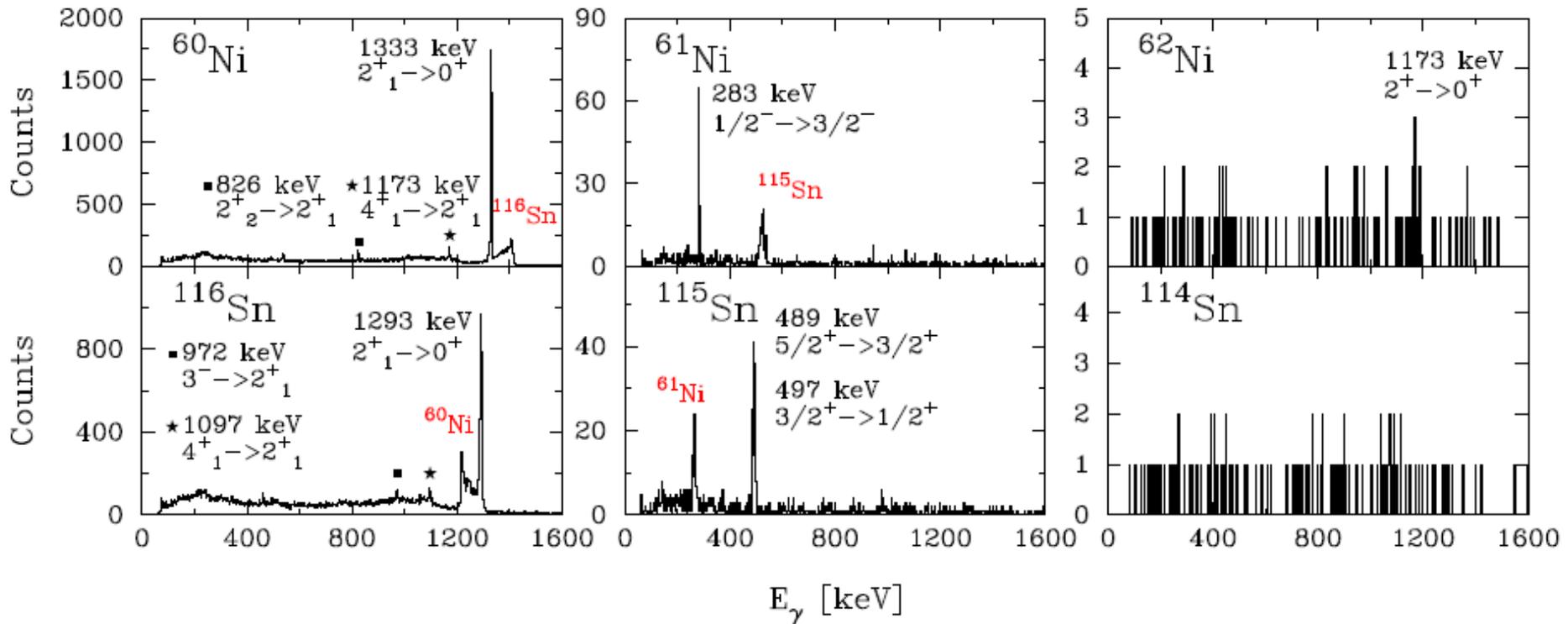
from the gamma array one gets the strength to excited states  $\longrightarrow \sigma_{exc}$

from the magnetic spectrometer one gets A,Z,Q (inclusive)  $\longrightarrow \sigma_{tot}$

simplifying...

$$\sigma_{g.s.} = \sigma_{tot} - \sigma_{exc}$$

**Pair neutron transfer probed via  $\gamma$ -particle coincidence  
in the  $^{60}\text{Ni}+^{116}\text{Sn}$  system at  $E_{\text{LAB}}=245\text{ MeV}$  and  $\theta_{\text{LAB}}=70^\circ$**



yields normalized to the  $2^+$  strength in  $^{60}\text{Ni}$

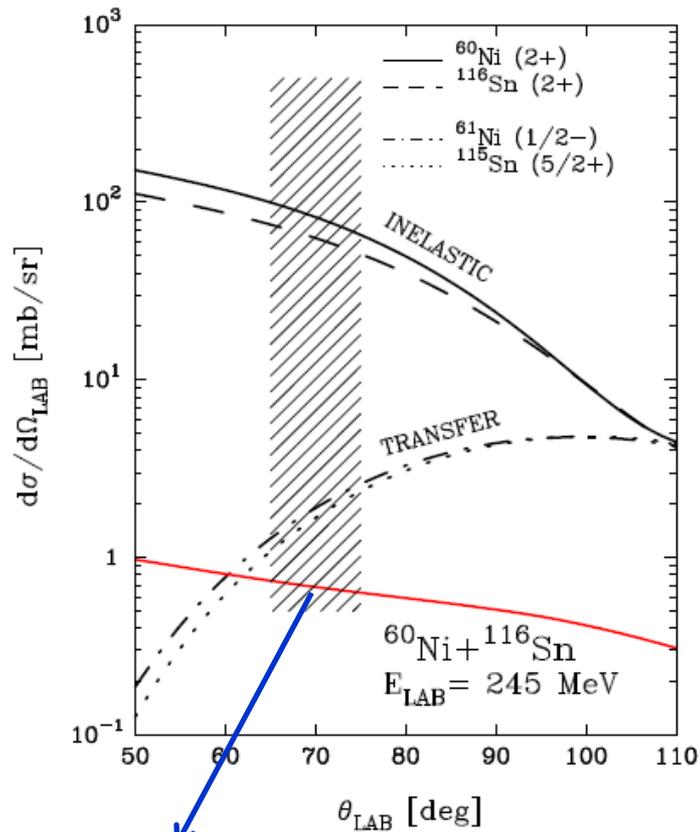
	Experiment	Theory
$^{116}\text{Sn}(2^+)$	$0.792 \pm 0.160$	0.720
$^{116}\text{Sn}(4_1^+)$	$0.042 \pm 0.011$	0.056
$^{60}\text{Ni}(4_1^+)$	$0.060 \pm 0.013$	0.11
$^{115}\text{Sn}(5/2^+)$	$0.018 \pm 0.003$	0.037
$^{61}\text{Ni}(1/2^-)$	$0.014 \pm 0.003$	0.033
$^{62}\text{Ni}(2^+)$	$< 0.00145$	-

**Experiment performed with  
PRISMA coupled to the AGATA  
DEMONSTRATOR**

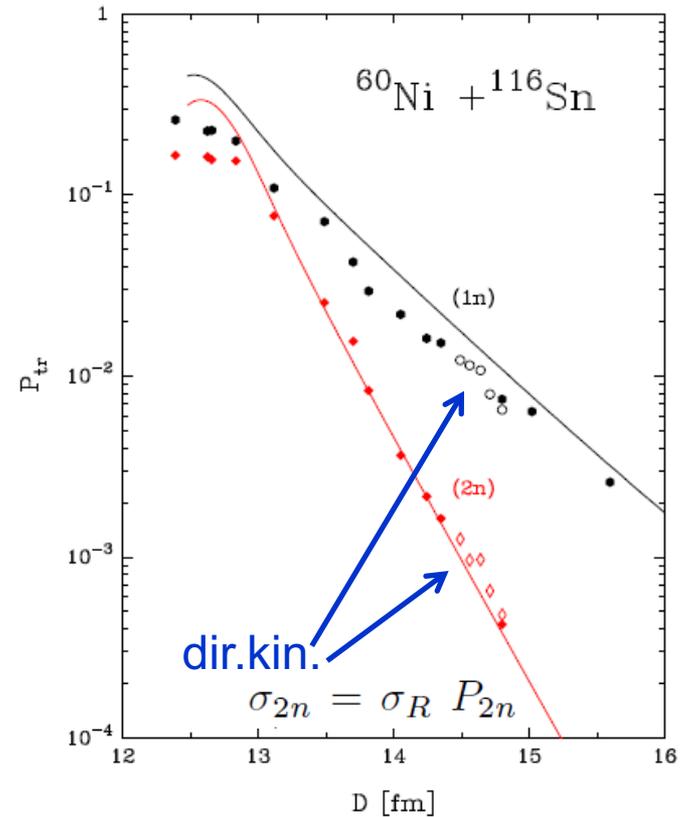
**D.Montanari, L.Corradi, S.Szilner,  
G.Pollarolo et al, PRC93(2016)054623**

# Pair neutron transfer probed via $\gamma$ -particle coincidence

by using the strength obtained with the gamma data for the inelastic and neutron transfer channels and with information from coupled-channel calculations we were able to quote that the fraction of the 2n channel populating the ground to ground state is larger than 76 %



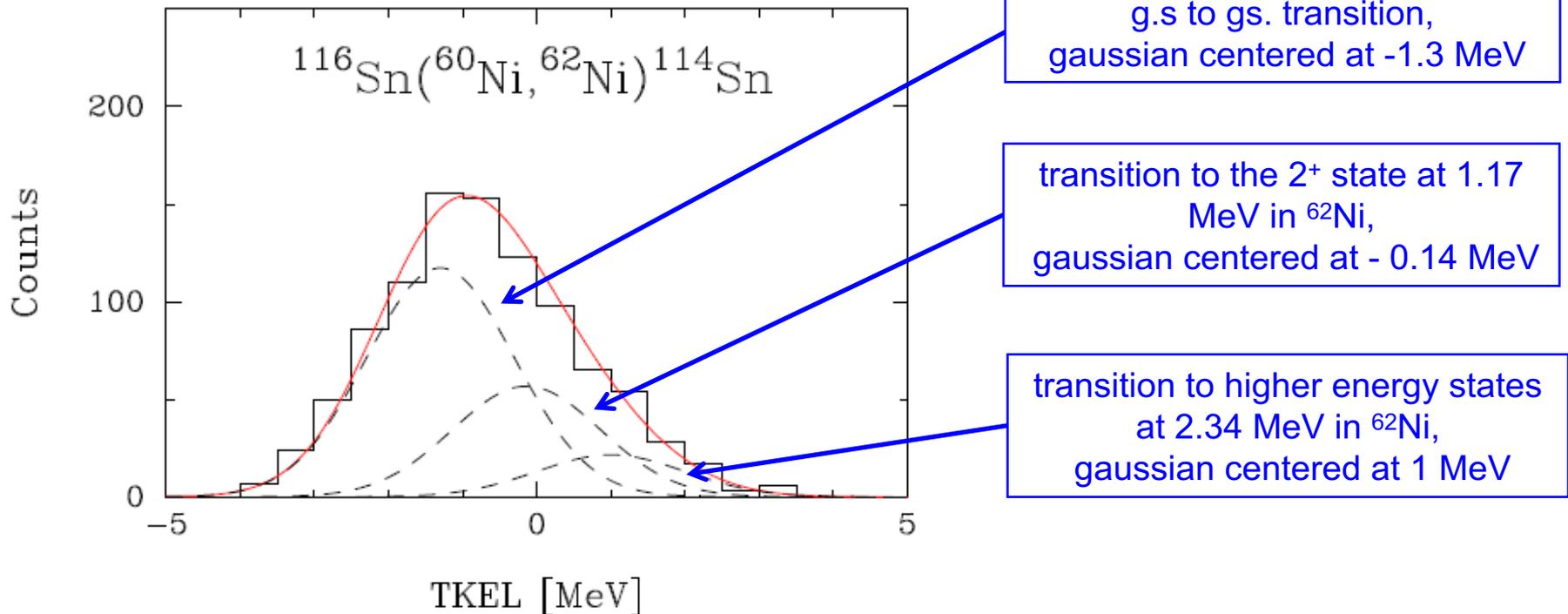
$$\sigma_R \left( 1 - \frac{\sigma_{\text{el}}}{\sigma_R} \right) = \sigma(2^+, ^{60}\text{Ni}) \left( 1 + \frac{\sigma(2^+, ^{116}\text{Sn})}{\sigma(2^+, ^{60}\text{Ni})} \right)$$



consistent comparison between inverse kinematics and direct kinematics data

# Pair neutron transfer probed via $\gamma$ -particle coincidence

## Analysis of the experimental TKEL distributions measured with PRISMA



**the analysis shows that the  $2^+$  states contribute 29% to the total transfer strength, in reasonable good agreement with the 24% estimation from the gamma strength**

**Sub-barrier transfer : much more work to be done**

**gamma-particle coincidences**

**proton transfer channels at large D**

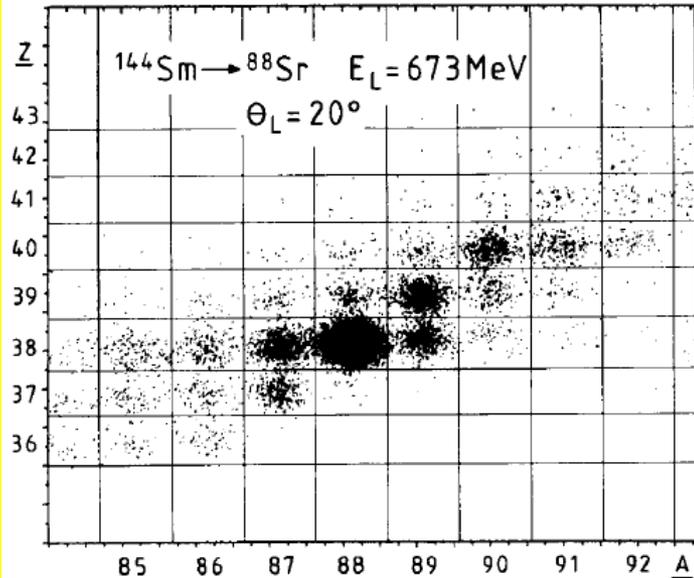
**very heavy systems**

**reactions with proton rich nuclei (also for np correlations)**

**reactions with neutron rich nuclei (density dependent forces)**

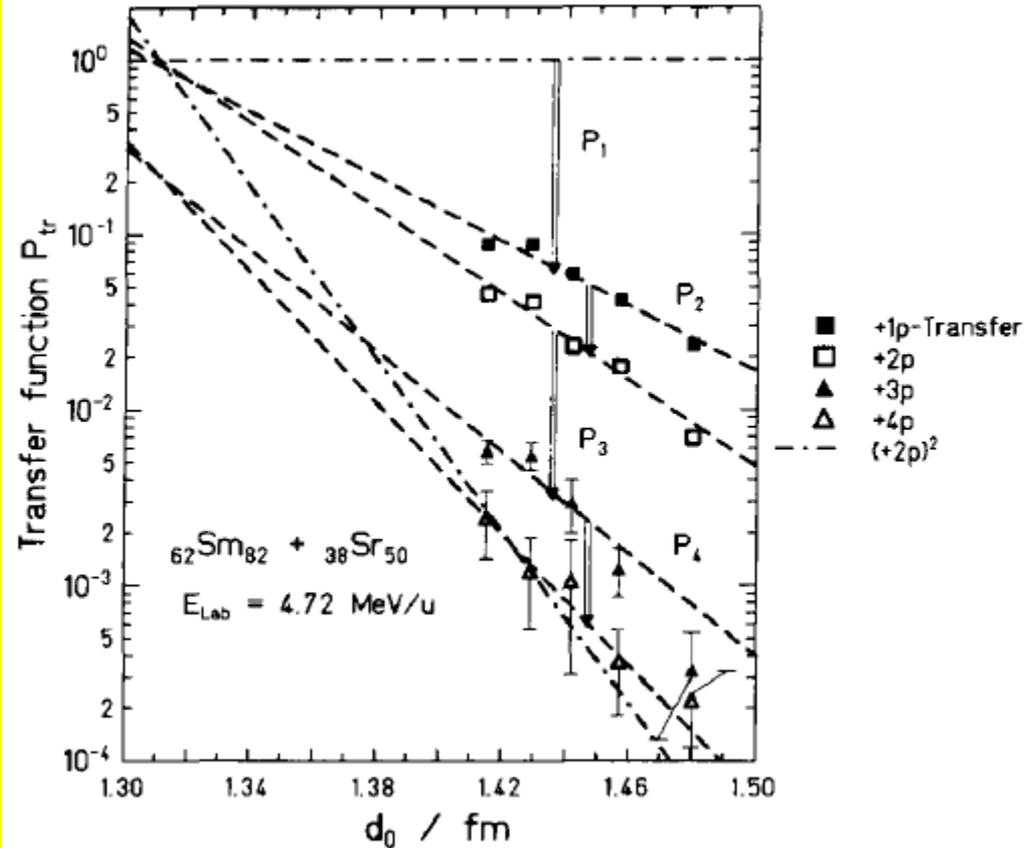
**need to develop calculations incorporating nn correlations,  
in general for high multipolarity states**

## Sub-barrier transfer : proton channels

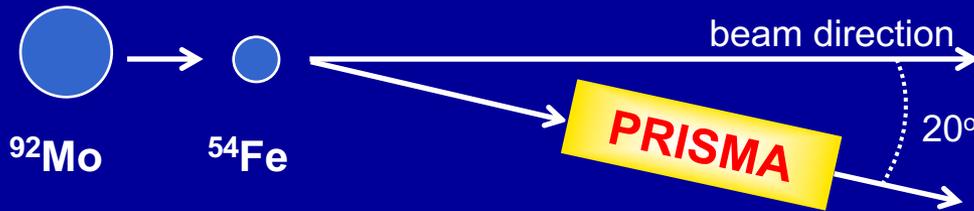


data are available, but for small  $D$ 's, where absorption plays a big role

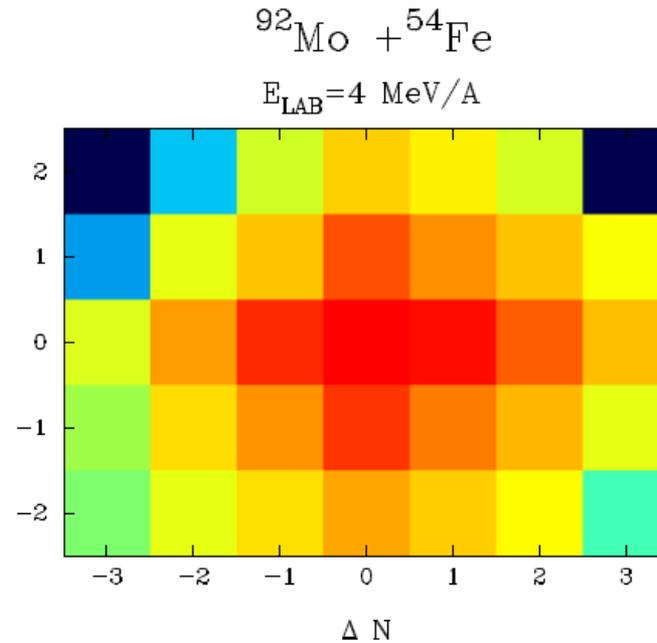
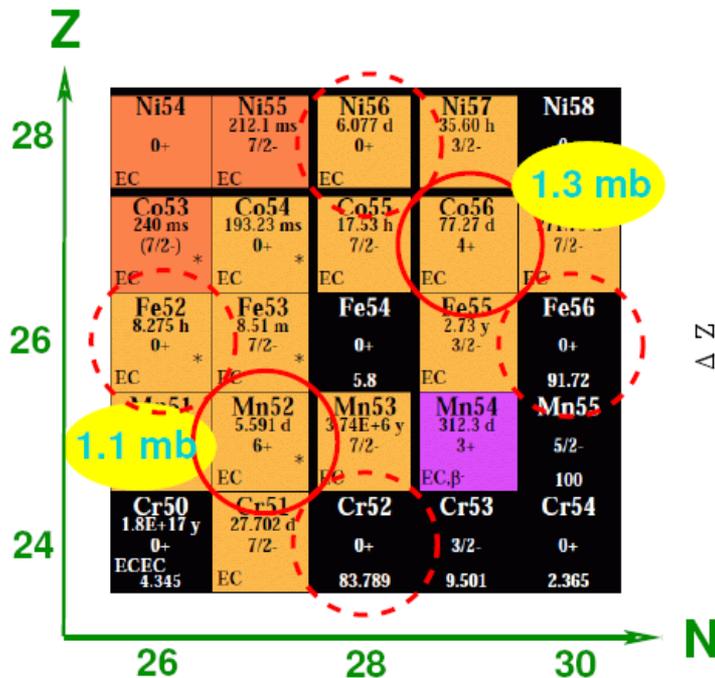
the interpretation of the enhancement factors is done at the phenomenological level



# Sub-barrier transfer : np correlations in proton-rich nuclei



Study of pair-correlation properties, populating at once  $\pm(nn)$ ,  $\pm(pp)$  and  $\pm(np)$  close to the  $N = Z = 27$  region



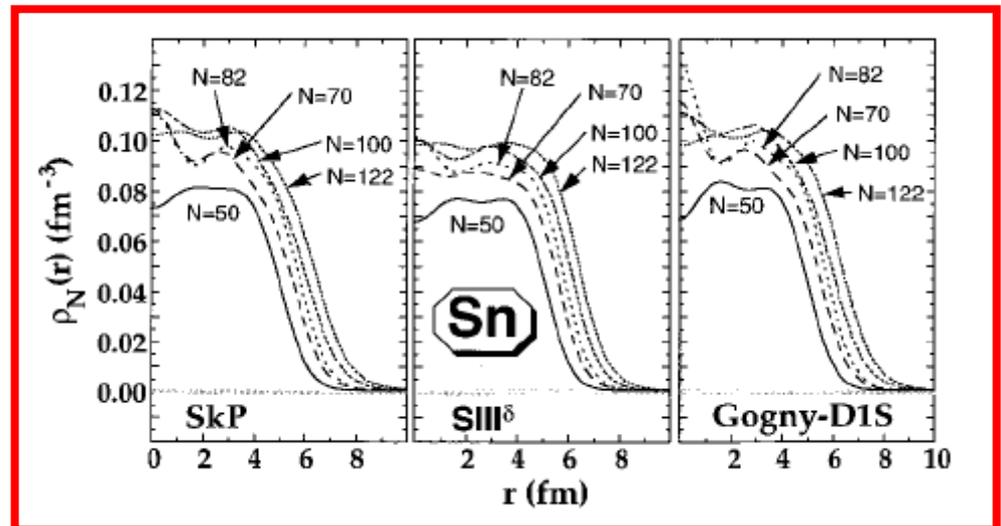
## Sub-barrier transfer : neutron rich nuclei

with RIB's, through the behaviour of transfer probabilities we may be able to probe, for instance

- onset of density dependent forces

$$V_{eff} = \delta(\vec{r}_1 - \vec{r}_2) \left( v_0 + v_p \left( \frac{\rho((\vec{r}_1 + \vec{r}_2)/2)}{\rho_c} \right)^p \right)$$

- neutron density profile





**L.Corradi, E.Fioretto, A.M.Stefanini,  
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**G.Pollarolo**

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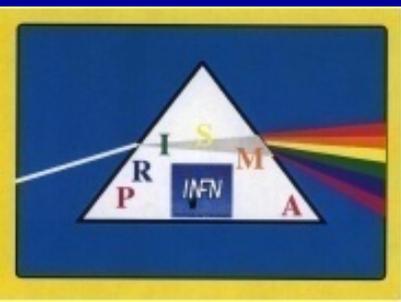
**S.Szilner, D.Jelavec-Malenica, T.Mijatovic,**

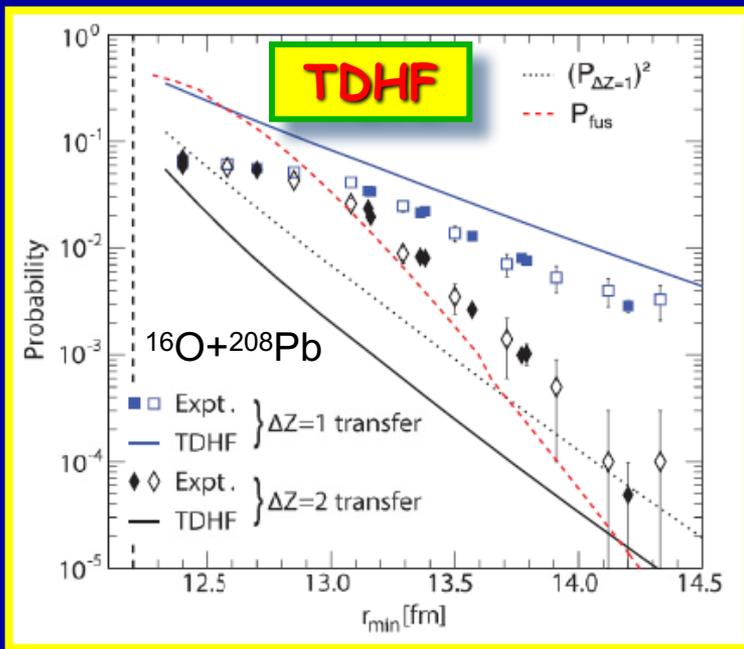
**N.Soic, M.Varga Pajtler**

*Ruđer Bošković Institute, Zagreb, Croatia*

**S.Courtin, A.Goasduff, F.Haas**

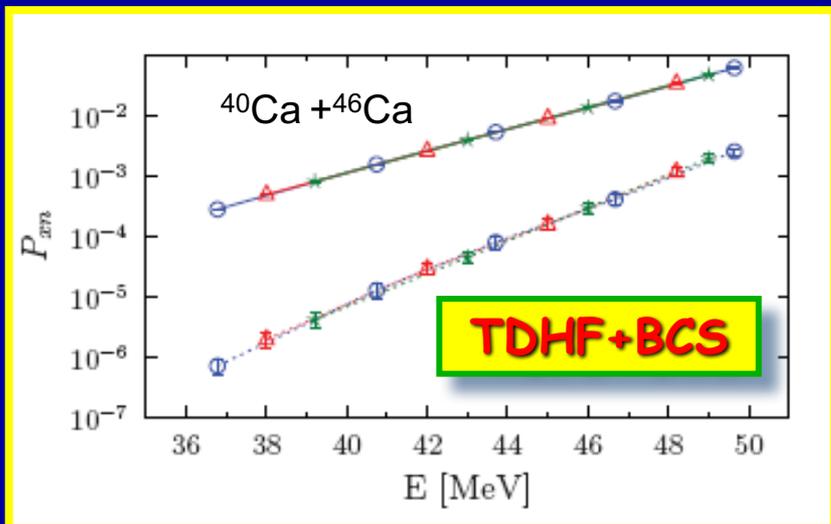
*IPHC, Strasbourg, France*



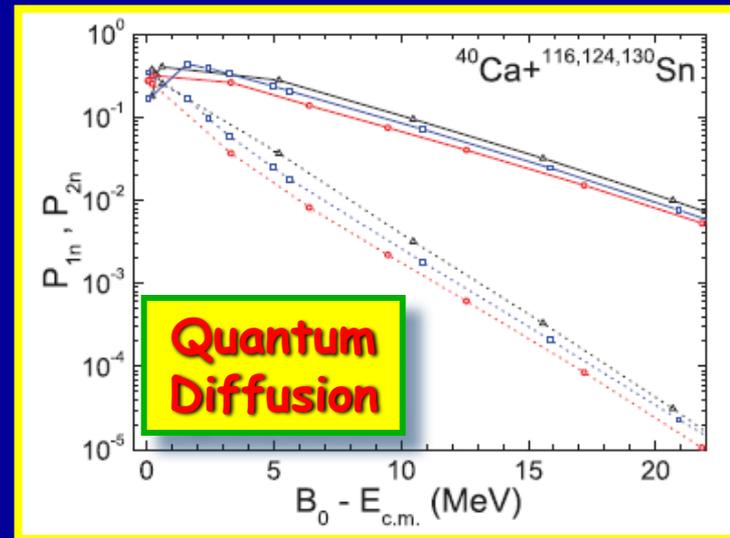


Sub barrier transfer reactions with heavy ions: new exp data are attracting a renewed interest in the theoretical community

C.Simene, PRL105(2010)192701  
 M.Evers et al, PRC84(2011)054614

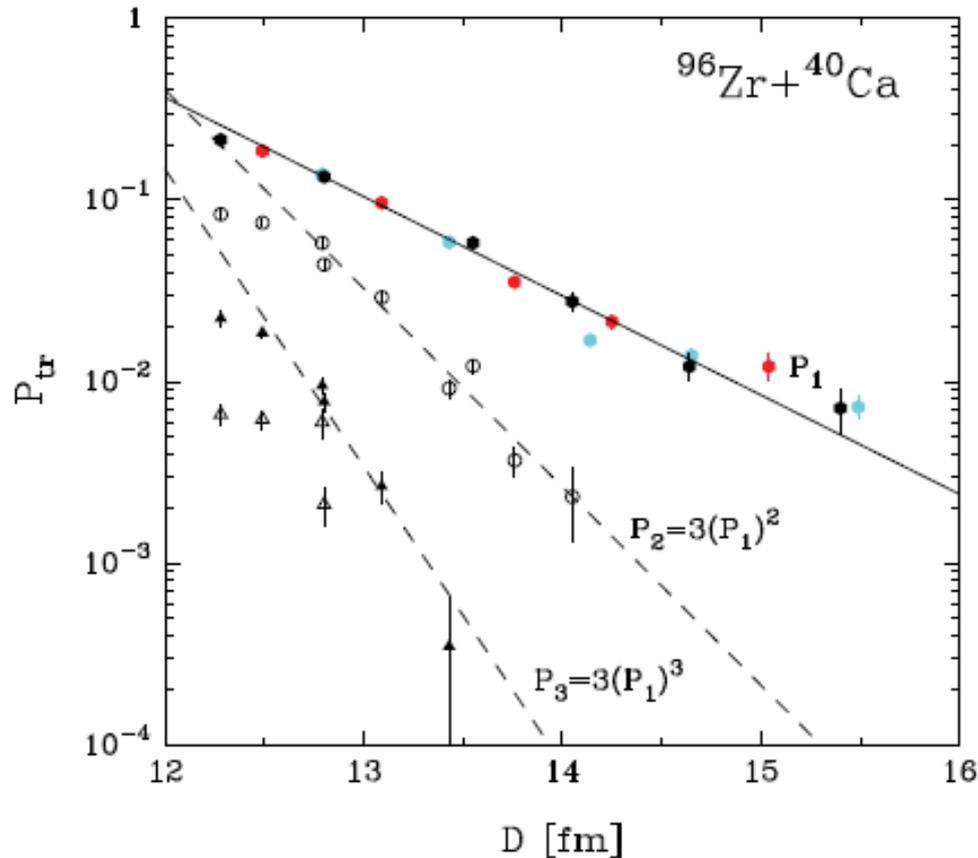


G.Scamps and D.Lacroix, PRC87(2013)014605



V.V.Sargsyan et al, PRC88(2013)064601

## Experimental transfer probabilities



$P_{tr}$  slope

$$P_{tr} \propto e^{-2\alpha D} \quad \alpha = \sqrt{\frac{2mB}{\hbar^2}}$$

$B \rightarrow$  binding energy

slopes of  $P_{tr}$  vs  $D$  are as expected from the binding energies (tail of the formfactor)

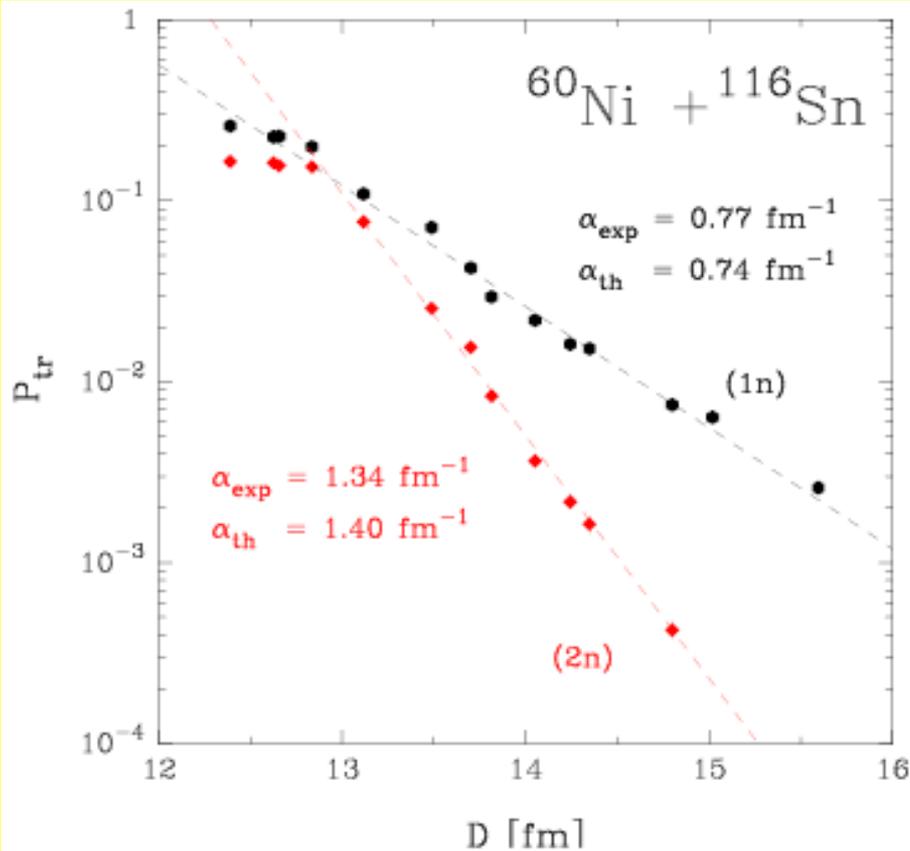
a bare phenomenological analysis shows an “enhanced” pair transfer,  $P_{2n} \sim 3(P_{1n})^2$  and  $P_{3n} \sim P_{1n}(P_{2n}) \sim 3(P_{1n})^3$

## $P_{tr}$ slope

$$P_{tr} \propto e^{-2\alpha D} \quad \alpha = \sqrt{\frac{2mB}{\hbar^2}}$$

B → binding energy

The  $^{60}\text{Ni} + ^{116}\text{Sn}$  system :  
experimental results

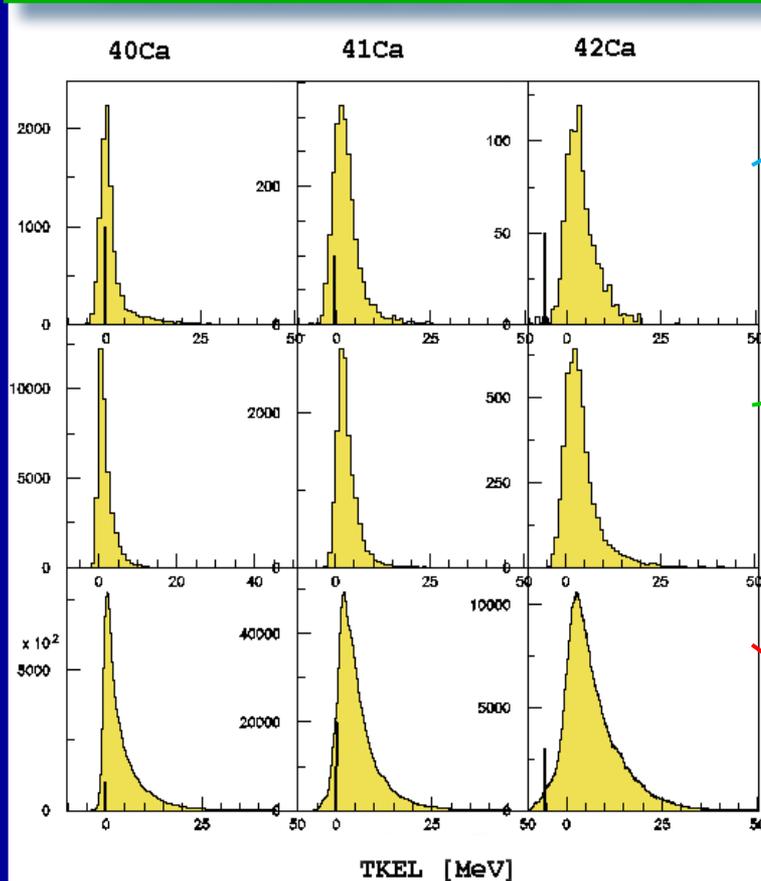


experimental slopes well  
match the ones expected  
from the binding energies

one gets an enhancement  
factor  $\sim 5.5$ , i.e.  
significantly larger than  
usual for pure neutron  
transfer

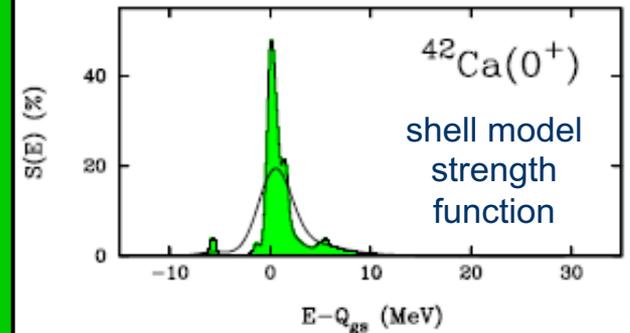
# Sub-barrier transfer : gamma-particle coincidences

## Total Kinetic Energy Loss distributions in $^{40}\text{Ca}+^{208}\text{Pb}$

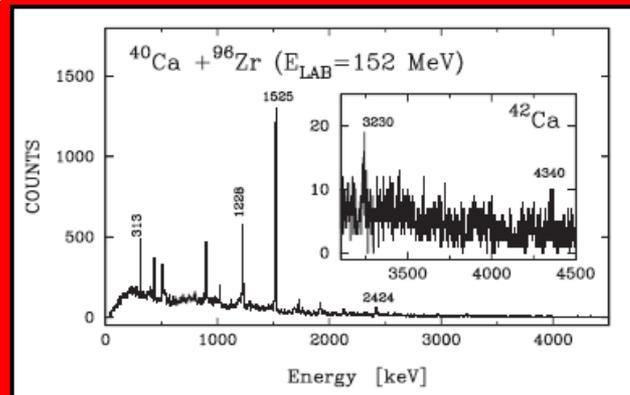


strong population close to the pairing vibrational region in  $^{42}\text{Ca}$

$^{96}\text{Zr}+^{40}\text{Ca}$  330 MeV - PRISMA  
L.Corradi et al. PRC84(2011)034603

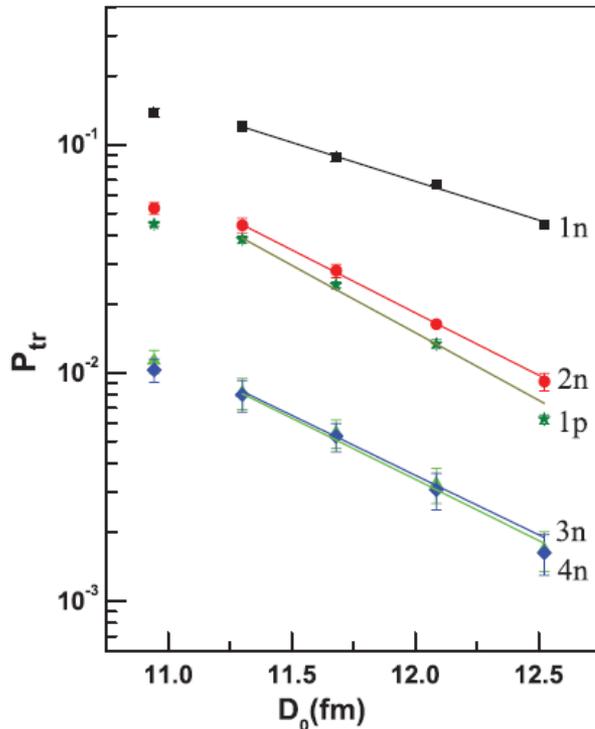
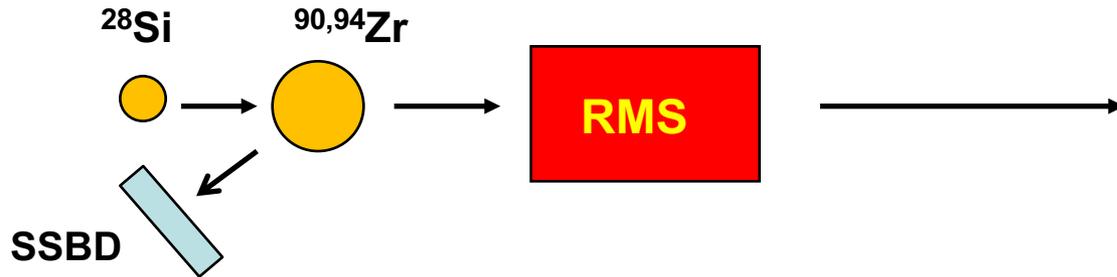


$^{40}\text{Ca}+^{208}\text{Pb}$  236 MeV - PISOLO  
S.Szilner et al PRC71(2005)044610



$^{40}\text{Ca}+^{96}\text{Zr}$  152 MeV - PRISMA+CLARA  
S.Szilner et al, PRC7(2007)024604

# Detection of (heavy) target like ions with recoil mass spectrometers



$$\alpha = \sqrt{\frac{2\mu B}{\hbar^2}}$$

$$\frac{P_{tr}}{\sin(\theta_{c.m.}/2)} \propto \exp(-2\alpha D)$$

data collected for multinucleon transfer in the range 11-12.5 fm

slopes much lower than those expected from binding energies

